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# Vibration response for a multilayer cylinder

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# Vibration response for a multilayer cylinder

David H. Chambers

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## 1 Analytical solution

Consider a circular cylinder composed of  $N$  concentric layers of materials. Each layer is characterized by its material density  $\rho_n$ , and elastic moduli  $\lambda_n$  and  $\mu_n$  ( $n = 1, \dots, N$ ). The interfaces between layers are specified by the radii  $R_0, R_1, \dots, R_N$ , with  $R_0$  the inner surface and  $R_N$  the outer surface of the cylinder. For each layer, the radial displacement field  $u_n(r, \theta)$  and tangential displacement field  $v_n(r, \theta)$  can be expressed using two elastic potentials  $\phi_n(r, \theta)$  and  $\psi_n(r, \theta)$ :

$$\begin{aligned} u_n(r, \theta) &= \frac{\partial \phi_n}{\partial r} + \frac{1}{r} \frac{\partial \psi_n}{\partial \theta}, \\ v_n(r, \theta) &= \frac{1}{r} \frac{\partial \phi_n}{\partial \theta} - \frac{\partial \psi_n}{\partial r}. \end{aligned} \quad (1)$$

The radial normal stress  $\sigma_{rr}^{(n)}(r, \theta)$  and shear stress  $\sigma_{r\theta}^{(n)}(r, \theta)$  are given by

$$\begin{aligned} \sigma_{rr}^{(n)}(r, \theta) &= (\lambda_n + 2\mu_n) \frac{\partial^2 \phi_n}{\partial r^2} + \frac{\lambda_n}{r} \left( \frac{\partial \phi_n}{\partial r} + \frac{1}{r} \frac{\partial^2 \phi_n}{\partial \theta^2} \right) + \frac{2\mu_n}{r} \left( \frac{\partial}{\partial r} - \frac{1}{r} \right) \frac{\partial \psi_n}{\partial \theta}, \\ \sigma_{r\theta}^{(n)}(r, \theta) &= \frac{2\mu_n}{r} \left( \frac{\partial}{\partial r} - \frac{1}{r} \right) \frac{\partial \phi_n}{\partial \theta} + \mu_n \left( -\frac{\partial^2 \psi_n}{\partial r^2} + \frac{1}{r} \frac{\partial \psi_n}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \psi_n}{\partial \theta^2} \right). \end{aligned} \quad (2)$$

The potentials are solutions to two Helmholtz equations:

$$\nabla^2 \phi_n + k_{Ln}^2 \phi_n = 0, \quad k_{Ln}^2 = \frac{\omega^2}{c_{Ln}^2} = \frac{\rho_n \omega^2}{\lambda_n + 2\mu_n}; \quad (3)$$

$$\nabla^2 \psi_n + k_{Tn}^2 \psi_n = 0, \quad k_{Tn}^2 = \frac{\omega^2}{c_{Tn}^2} = \frac{\rho_n \omega^2}{\mu_n}. \quad (4)$$

For cylindrical geometry, the solutions to the Helmholtz equations can be written as Fourier series in the angular coordinate  $\theta$ :

$$\phi_n(r, \theta) = \sum_{m=0}^{\infty} \varepsilon_m (-1)^m \left[ A_n^{(m)} J_m(k_{Ln} r) + B_n^{(m)} Y_m(k_{Ln} r) \right] \cos(m\theta), \quad (5)$$

$$\psi_n(r, \theta) = \sum_{m=0}^{\infty} \varepsilon_m (-1)^m \left[ D_n^{(m)} J_m(k_{Tn} r) + E_n^{(m)} Y_m(k_{Tn} r) \right] \sin(m\theta), \quad (6)$$

where

$$\varepsilon_m = \begin{cases} 1, & m = 0 \\ 2, & m \geq 1 \end{cases}. \quad (7)$$

Equations for the coefficients  $A_n^{(m)}$ ,  $B_n^{(m)}$ ,  $D_n^{(m)}$ , and  $E_n^{(m)}$ , are obtained by requiring continuity of displacement and stress across each interface ( $r = R_1, R_2, \dots, R_{N-1}$ ) and by specifying the tractions on the inner

( $r = R_0$ ) and outer ( $r = R_N$ ) surfaces. These conditions can be applied independently for each order  $m$  in the Fourier series. The final matrix equation for the coefficients of the  $m$ th mode is  $\mathbf{Z}^{(m)} \cdot \mathbf{C}^{(m)} = \mathbf{b}^{(m)}$ , where  $\mathbf{C}^{(m)}$  is a vector containing the coefficients,  $\mathbf{b}^{(m)}$  is a vector determined from the applied tractions, and  $\mathbf{Z}^{(m)}$  is a matrix consisting of Bessel functions evaluated at the interfaces.

The vector  $\mathbf{C}^{(m)}$  can be constructed by stacking  $N$  vectors  $\mathbf{C}_n^{(m)}$ :

$$\mathbf{C}^{(m)} = \begin{bmatrix} \mathbf{C}_1^{(m)} \\ \mathbf{C}_2^{(m)} \\ \vdots \\ \mathbf{C}_N^{(m)} \end{bmatrix}, \quad \mathbf{C}_n^{(m)} = \begin{bmatrix} A_n^{(m)} \\ B_n^{(m)} \\ D_n^{(m)} \\ E_n^{(m)} \end{bmatrix}. \quad (8)$$

The matrix  $\mathbf{Z}^{(m)}$  can be written in terms of the  $4 \times 4$  matrices  $\mathbf{S}_n^{(m)}(r)$  and the  $2 \times 4$  matrices  $\mathbf{F}_n^{(m)}(r)$ :

$$\mathbf{Z}^{(m)} = \begin{bmatrix} \mathbf{F}_1^{(m)}(R_0) & \mathbf{0} & \dots & \mathbf{0} \\ \mathbf{S}_1^{(m)}(R_1) & -\mathbf{S}_2^{(m)}(R_1) & \ddots & \vdots \\ \mathbf{0} & \ddots & \ddots & \mathbf{0} \\ \vdots & \ddots & \mathbf{S}_{N-1}^{(m)}(R_{N-1}) & -\mathbf{S}_N^{(m)}(R_{N-1}) \\ \mathbf{0} & \dots & \mathbf{0} & \mathbf{F}_N^{(m)}(R_N) \end{bmatrix}. \quad (9)$$

The  $\mathbf{S}_n^{(m)}(r)$  matrices are composed of the  $\mathbf{F}_n^{(m)}(r)$  matrices and another  $2 \times 4$  matrix  $\mathbf{G}_n^{(m)}(r)$ :

$$\mathbf{S}_n^{(m)}(r) = \begin{bmatrix} \mathbf{F}_n^{(m)}(r) \\ \mathbf{G}_n^{(m)}(r) \end{bmatrix}, \quad (10)$$

where

$$\mathbf{G}_n^{(m)}(r) = \begin{bmatrix} k_{Ln} J'_m(k_{Ln}r) & k_{Ln} Y'_m(k_{Ln}r) & \frac{m}{r} J_m(k_{Tn}r) & \frac{m}{r} Y_m(k_{Tn}r) \\ \frac{m}{r} J_m(k_{Ln}r) & \frac{m}{r} Y_m(k_{Ln}r) & k_{Tn} J'_m(k_{Tn}r) & k_{Tn} Y'_m(k_{Tn}r) \end{bmatrix}. \quad (11)$$

Finally, the  $\mathbf{F}_n^{(m)}(r)$  matrix is given by

$$\mathbf{F}_n^{(m)}(r) = \begin{bmatrix} F_{11n}^{(m)}(k_{Ln}r) & F_{12n}^{(m)}(k_{Ln}r) & \frac{2\rho_n \omega c_{Tn} m}{r} g_m^{(1)}(k_{Tn}r) & \frac{2\rho_n \omega c_{Tn} m}{r} g_m^{(2)}(k_{Tn}r) \\ \frac{2\rho_n \omega c_{Tn}^2 m}{c_{Ln} r} g_m^{(1)}(k_{Ln}r) & \frac{2\rho_n \omega c_{Tn}^2 m}{c_{Ln} r} g_m^{(2)}(k_{Ln}r) & F_{23n}^{(m)}(k_{Tn}r) & F_{24n}^{(m)}(k_{Tn}r) \end{bmatrix}, \quad (12)$$

with

$$\begin{bmatrix} F_{11n}^{(m)}(k_{Ln}r) & F_{12n}^{(m)}(k_{Ln}r) \end{bmatrix} = -\rho_n \omega^2 [J_m(k_{Ln}r) \quad Y_m(k_{Ln}r)] - \frac{2\rho_n \omega c_{Tn}^2}{c_{Ln} r} [f_m^{(1)}(k_{Ln}r) \quad f_m^{(2)}(k_{Ln}r)], \quad (13)$$

and

$$\begin{bmatrix} F_{23n}^{(m)}(k_{Tn}r) & F_{24n}^{(m)}(k_{Tn}r) \end{bmatrix} = \rho_n \omega^2 [J_m(k_{Tn}r) \quad Y_m(k_{Tn}r)] + \frac{2\rho_n \omega c_{Tn}}{r} [f_m^{(1)}(k_{Tn}r) \quad f_m^{(2)}(k_{Tn}r)]. \quad (14)$$

The functions  $f_m^{(1)}$ ,  $f_m^{(2)}$ ,  $g_m^{(1)}$ , and  $g_m^{(2)}$  are combinations of Bessel functions given below:

$$\begin{aligned} f_m^{(1)}(x) &= J'_m(x) - \frac{m^2}{x} J_m(x) \\ f_m^{(2)}(x) &= Y'_m(x) - \frac{m^2}{x} Y_m(x) \\ g_m^{(1)}(x) &= J'_m(x) - \frac{1}{x} J_m(x) \\ g_m^{(2)}(x) &= J'_m(x) - \frac{1}{x} J_m(x) \end{aligned} \quad (15)$$

This completes the specification of the matrix  $\mathbf{Z}^{(m)}$ .

The last part of the solution is determining the forcing vector  $\mathbf{b}^{(m)}$  for each mode. We will consider two cases, normal point load on the inside surface  $r = R_0$ , and normal point load on the outer surface  $r = R_N$ . In either case we can assume the load is applied at angular position  $\theta = 0$  with no loss of generality. The mathematical expression for the normal point load is

$$\sigma_{rr}(R, \theta) = -\frac{1}{R}\delta(\theta) = -\frac{1}{2\pi R} \sum_{m=0}^{\infty} \varepsilon_m \cos(m\theta), \quad (16)$$

with  $R = R_0$  for the inside surface, or  $R = R_N$  for the outside surface. After substituting the Fourier series (5) into the expression for the normal stress in (2), then equating it with (16) above, we obtain the following for the forcing vectors  $\mathbf{b}_0^{(m)}$  and  $\mathbf{b}_N^{(m)}$  for the inside and outside surfaces:

$$\mathbf{b}_0^{(m)} = -\frac{i^m}{2\pi R_0} \begin{bmatrix} 1 \\ 0 \\ \vdots \\ 0 \end{bmatrix}, \quad \mathbf{b}_N^{(m)} = -\frac{i^m}{2\pi R_N} \begin{bmatrix} 0 \\ \vdots \\ 0 \\ 1 \\ 0 \end{bmatrix}. \quad (17)$$

Once the coefficients for each mode are obtained by solving  $\mathbf{Z}^{(m)} \cdot \mathbf{C}^{(m)} = \mathbf{b}^{(m)}$ , we would like to calculate the normal displacement on the inner and outer surfaces. Using the series expansions (5) for the potentials and their relationships to the displacements (1) we obtain the following:

$$u_0(\theta) = u(R_0, \theta) = k_{L0} \sum_{m=0}^{\infty} \varepsilon_m (-i)^m \left\{ J'_m(k_{L0}R_0) A_0^{(m)} + Y'_m(k_{L0}R_0) B_0^{(m)} + \frac{m}{k_{L0}R_0} \left[ J_m(k_{T0}R_0) D_0^{(m)} + \frac{m}{R_0} Y_m(k_{T0}R_0) E_0^{(m)} \right] \right\}, \quad (18)$$

$$u_N(\theta) = u(R_N, \theta) = k_{LN} \sum_{m=0}^{\infty} \varepsilon_m (-i)^m \left\{ J'_m(k_{LN}R_N) A_N^{(m)} + Y'_m(k_{LN}R_N) B_N^{(m)} + \frac{m}{k_{LN}R_N} \left[ J_m(k_{TN}R_N) D_N^{(m)} + \frac{m}{R_N} Y_m(k_{TN}R_N) E_N^{(m)} \right] \right\}. \quad (19)$$

Suppose we sample the normal displacements at angular positions  $\theta = \theta_n = 2\pi n/N$  for  $n = 1, \dots, N-1$ , then we can extract the  $j$ th term ( $j \leq N$ ) of either (18) or (19) by evaluating the weighted sums

$$u_{0j} = \sum_{n=0}^{N-1} u_0(\theta_n) \cos\left(\frac{2\pi j n}{N}\right) \quad (20)$$

for the inner surface, and

$$u_{Nj} = \sum_{n=0}^{N-1} u_N(\theta_n) \cos\left(\frac{2\pi j n}{N}\right) \quad (21)$$

for the outer surface. These are the quantities are plotted in the text to determine the vibrational response in each example.

Finally, there are two special cases that can occur: a fluid layer ( $\mu_n = 0$  for some  $n$ ), and a solid core cylinder ( $R_0 = 0$ ). In a fluid the shear stress is zero, which eliminates the potential  $\psi_n$ , so that  $D_n^{(m)} = E_n^{(m)} = 0$  identically. In addition, the shear stress in adjacent layers must go to zero at the interfaces with the fluid layer ( $\sigma_{r\theta}^{(n-1)}(R_{n-1}, \theta) = 0$  and  $\sigma_{r\theta}^{(n+1)}(R_n, \theta) = 0$ ). Similarly, there is no continuity condition for the tangential displacements ( $v_{n-1}$  and  $v_{n+1}$ ) in adjacent layers. Thus we can eliminate the columns of

$Z^{(m)}$  corresponding to  $D_n^{(m)}$  and  $E_n^{(m)}$ , along with the rows of  $Z^{(m)}$  containing the first and third rows of  $\mathbf{S}_n^{(m)}$ . This modification can be done for each fluid layer to determine the final form for  $Z^{(m)}$ .

If the cylinder has a solid core ( $R_0 = 0$ ), then the condition that the potentials must be finite at  $r = 0$  forces  $B_1^{(m)} = E_1^{(m)} = 0$ . It also eliminates the free surface condition for the innermost layer. We can then eliminate the first two rows of  $Z^{(m)}$  along with the second and fourth columns which correspond to  $B_1^{(m)}$  and  $E_1^{(m)}$  to determine the final form for  $Z^{(m)}$ .

## 2 Acknowledgements / Disclaimer

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```

function [ampin,ampout] = measure_cyl(layer_parms,w,Z,maxorder,src_vec)
% function [ampin,ampout] = measure_cyl(layer_parms,w,Z,maxorder,src_vec)
% This function calculates the "measured" modal amplitudes
% for a cylinder represented by the matrix Z, due to a
% source. Inputs are:
%   layer_parms:      structure of layer parameters
%     .R:            interface radii
%     .mu:           shear modulus (Lame)
%     .lambda:        compression modulus (Lame)
%     .rho:           density
%   w:             angular frequency
%   Z:             matrix for modal coefficients
%   maxorder:       maximum order of angular dependence (0:maxorder)
%   src_vec:        source vector
% Note ampin = 0 if cylinder is solid (R(1)=0).
% Created 7/11/06

```

```

R = layer_parms.R;
mu = layer_parms.mu;
imu0 = find(abs(mu)==0);
imu0c = find(abs(mu)>0);
lambda = layer_parms.lambda;
rho = layer_parms.rho;
nlay = length(rho);
nZ0 = 4*nlay;
nR = nlay+1;
cL = sqrt((lambda+2*mu)./rho);
cT = sqrt(mu./rho);
kL = w./cL;
kT = zeros(nlay,1);
if length(imu0)>0
    kT(imu0) = 1.e8;
    if length(imu0c)>0
        kT(imu0c) = w./cT(imu0c);
    end
else
    kT = w./cT;
end
muflag0 = sign(abs(mu));
m = (0:maxorder)';
mmax = maxorder+1;
JL = zeros(mmax,1);
JT = zeros(mmax,1);
YL = zeros(mmax,1);
YT = zeros(mmax,1);
JpL = zeros(mmax,1);
YpL = zeros(mmax,1);
JpT = zeros(mmax,1);
YpT = zeros(mmax,1);
gi = zeros(mmax,4);
go = zeros(mmax,4);
ampin = zeros(mmax,1);
ampout = zeros(mmax,1);

argLo = kL(nlay)*R(nR);
argTo = kT(nlay)*R(nR);
[JLo,YLo,JpLo,YpLo] = make_bessel(maxorder,argLo,0);

```

```

[JTo,YTo,JpTo,YpTo] = make_bessel(maxorder,argTo,0);
go(:,1:2) = kL(nlay)*[JpLo YpLo];
go(:,3) = muflag0(nlay)*m.*JTo/R(nR);
go(:,4) = muflag0(nlay)*m.*YTo/R(nR);
if R(1)>0
    argLi = kL(1)*R(1);
    argTi = kT(1)*R(1);
    [JLi,YLi,JpLi,YpLi] = make_bessel(maxorder,argLi,0);
    [JTi,YTi,JpTi,YpTi] = make_bessel(maxorder,argTi,0);
    gi(:,1:2) = kL(1)*[JpLi YpLi];
    gi(:,3) = muflag0(1)*m.*JTi/R(1);
    gi(:,4) = muflag0(1)*m.*YTi/R(1);
end

for jm = 1:mmax
    tmp = zeros(nZ0,1);
    Z1 = squeeze(Z(:,:,jm));
    irZ1 = sign(sum(abs(Z1),2));
    itmp = find(irZ1);
    icZ1 = sign(sum(abs(Z1),1));
    nZ = sum(irZ1);
    if nZ<nZ0
        Z2 = zeros(nZ,nZ);
        indZ = find(irZ1*icZ1>0);
        Z2 = Z1(indZ);
        Z2 = reshape(Z2,nZ,nZ);
        src_vec2 = src_vec(itmp);
        if rcond(Z2)<1.e-12
            tmp2 = pinv(Z2)*src_vec2;
        else
            tmp2 = Z2\src_vec2;
        end
        tmp(itmp) = tmp2;
    else
        if rcond(Z1)<1.e12
            tmp = pinv(Z1)*src_vec;
        else
            tmp = Z1\src_vec;
        end
    end
    ampout(jm) = squeeze(go(jm,:))*tmp(((4*nlay-3):(4*nlay)));
    if R(1)>0
        ampin(jm) = squeeze(gi(jm,:))*tmp(1:4);
    end
end

% end function

```

```

function Z = make_Z(maxorder,layer_parms,w)
% function Z = make_Z(maxorder,layer_parms,w)
% This function creates the Z matrix for a single frequency from the
% layer parameters. Input variables are:
% maxorder: maximum order of angular dependence (0:maxorder)
% layer_parms: structure of layer parameters
%     .R:           interface radii
%     .mu:          shear modulus (Lame)
%     .lambda:      compression modulus (Lame)
%     .rho:          density
% w:             angular frequency
% Created 7/10/06

R = layer_parms.R;
mu = layer_parms.mu;
imu0 = find(abs(mu)==0);
imu0c = find(abs(mu)>0);
lambda = layer_parms.lambda;
rho = layer_parms.rho;
nlay = length(rho);
nR = nlay+1;
cL = sqrt((lambda+2*mu)./rho);
cT = sqrt(mu./rho);
kL = w./cL;
kT = zeros(nlay,1);
if length(imu0)>0
    kT(imu0) = 1.e8;
    if length(imu0c)>0
        kT(imu0c) = w./cT(imu0c);
    end
else
    kT = w./cT;
end
muflag0 = sign(abs(mu));
m = (0:maxorder)';
mmax = maxorder+1;
JL = zeros(mmax,1);
JT = zeros(mmax,1);
YL = zeros(mmax,1);
YT = zeros(mmax,1);
JpL = zeros(mmax,1);
YpL = zeros(mmax,1);
JpT = zeros(mmax,1);
YpT = zeros(mmax,1);
F1L = zeros(mmax,1);
F2L = zeros(mmax,1);
G1L = zeros(mmax,1);
G2T = zeros(mmax,1);
Z = zeros(4*nlay,4*nlay,mmax);

argL = kL(1)*R(1);
argT = kT(1)*R(1);
[JL,YL,JpL,YpL] = make_bessel(maxorder,argL,0);
[JT,YT,JpT,YpT] = make_bessel(maxorder,argT,0);
[F1L,F2L,G1L,G2L] = make_FG(maxorder,JL,YL,JpL,YpL,argL,0);
[F1T,F2T,G1T,G2T] = make_FG(maxorder,JT,YT,JpT,YpT,argT,0);
a1 = rho(1)*w^2;

```

```

if R(1)==0
    a2 = 0;
    a3 = 0;
else
    a2 = 2*mu(1)*kL(1)/R(1);
    a3 = 2*muflag0(1)*rho(1)*w^2/(kT(1)*R(1));
end
Z(1,1,:) = reshape(-a1*JL-a2*F1L,1,1,mmax);
Z(1,2,:) = reshape(-a1*YL-a2*F2L,1,1,mmax);
Z(1,3,:) = reshape(a3*m.*G1T,1,1,mmax);
Z(1,4,:) = reshape(a3*m.*G2T,1,1,mmax);
Z(2,1,:) = reshape(-a2*m.*G1L,1,1,mmax);
Z(2,2,:) = reshape(-a2*m.*G2L,1,1,mmax);
Z(2,3,:) = muflag0(1)*reshape(a1*JT+a3*F1T,1,1,mmax);
Z(2,4,:) = muflag0(1)*reshape(a1*YT+a3*F2T,1,1,mmax);

for i=2:nlay
    irow = 4*(i-1)-1;
    icoll1 = 4*(i-2)+1;
    icoll2 = 4*(i-1)+1;
    muprod0 = muflag0(i-1)*muflag0(i);

    argL = kL(i-1)*R(i);
    argT = kT(i-1)*R(i);
    [JL,YL,JpL,YpL] = make_bessel(maxorder,argL,0);
    [JT,YT,JpT,YpT] = make_bessel(maxorder,argT,0);
    [F1L,F2L,G1L,G2L] = make_FG(maxorder,JL,YL,JpL,YpL,argL,0);
    [F1T,F2T,G1T,G2T] = make_FG(maxorder,JT,YT,JpT,YpT,argT,0);
    a1 = rho(i-1)*w^2;
    a2 = 2*mu(i-1)*kL(i-1)/R(i);
    a3 = 2*muflag0(i-1)*rho(i-1)*w^2/(kT(i-1)*R(i));
    Z(irow,icoll1,:) = reshape(-a1*JL-a2*F1L,1,1,mmax);
    Z(irow,icoll1+1,:) = reshape(-a1*YL-a2*F2L,1,1,mmax);
    Z(irow,icoll1+2,:) = reshape(a3*m.*G1T,1,1,mmax);
    Z(irow,icoll1+3,:) = reshape(a3*m.*G2T,1,1,mmax);
    Z(irow+1,icoll1,:) = reshape(-a2*m.*G1L,1,1,mmax);
    Z(irow+1,icoll1+1,:) = reshape(-a2*m.*G2L,1,1,mmax);
    Z(irow+1,icoll1+2,:) = muflag0(i-1)*reshape(a1*JT+a3*F1T,1,1,mmax);
    Z(irow+1,icoll1+3,:) = muflag0(i-1)*reshape(a1*YT+a3*F2T,1,1,mmax);
    Z(irow+2,icoll1,:) = reshape(kL(i-1)*JpL,1,1,mmax);
    Z(irow+2,icoll1+1,:) = reshape(kL(i-1)*YpL,1,1,mmax);
    Z(irow+2,icoll1+2,:) = reshape(m.*JT/R(i),1,1,mmax);
    Z(irow+2,icoll1+3,:) = reshape(m.*YT/R(i),1,1,mmax);
    Z(irow+3,icoll1,:) = muprod0*reshape(m.*JL/R(i),1,1,mmax);
    Z(irow+3,icoll1+1,:) = muprod0*reshape(m.*YL/R(i),1,1,mmax);
    Z(irow+3,icoll1+2,:) = muprod0*reshape(kT(i-1)*JpT,1,1,mmax);
    Z(irow+3,icoll1+3,:) = muprod0*reshape(kT(i-1)*YpT,1,1,mmax);

    argL = kL(i)*R(i);
    argT = kT(i)*R(i);
    [JL,YL,JpL,YpL] = make_bessel(maxorder,argL,0);
    [JT,YT,JpT,YpT] = make_bessel(maxorder,argT,0);
    [F1L,F2L,G1L,G2L] = make_FG(maxorder,JL,YL,JpL,YpL,argL,0);
    [F1T,F2T,G1T,G2T] = make_FG(maxorder,JT,YT,JpT,YpT,argT,0);
    a1 = rho(i)*w^2;
    a2 = 2*mu(i)*kL(i)/R(i);
    a3 = 2*muflag0(i)*rho(i)*w^2/(kT(i)*R(i));

```

```

Z(irow,icol2,:) = -reshape(-a1*JL-a2*F1L,1,1,mmax);
Z(irow,icol2+1,:) = -reshape(-a1*YL-a2*F2L,1,1,mmax);
Z(irow,icol2+2,:) = -reshape(a3*m.*G1T,1,1,mmax);
Z(irow,icol2+3,:) = -reshape(a3*m.*G2T,1,1,mmax);
Z(irow+1,icol2,:) = -reshape(-a2*m.*G1L,1,1,mmax);
Z(irow+1,icol2+1,:) = -reshape(-a2*m.*G2L,1,1,mmax);
Z(irow+1,icol2+2,:) = -muflag0(i)*reshape(a1*JT+a3*F1T,1,1,mmax);
Z(irow+1,icol2+3,:) = -muflag0(i)*reshape(a1*YT+a3*F2T,1,1,mmax);
Z(irow+2,icol2,:) = -reshape(kL(i)*JpL,1,1,mmax);
Z(irow+2,icol2+1,:) = -reshape(kL(i)*YpL,1,1,mmax);
Z(irow+2,icol2+2,:) = -reshape(m.*JT/R(i),1,1,mmax);
Z(irow+2,icol2+3,:) = -reshape(m.*YT/R(i),1,1,mmax);
Z(irow+3,icol2,:) = -muprod0*reshape(m.*JL/R(i),1,1,mmax);
Z(irow+3,icol2+1,:) = -muprod0*reshape(m.*YL/R(i),1,1,mmax);
Z(irow+3,icol2+2,:) = -muprod0*reshape(kT(i)*JpT,1,1,mmax);
Z(irow+3,icol2+3,:) = -muprod0*reshape(kT(i)*YpT,1,1,mmax);
end

irow = 4*nlay-1;
icol = 4*nlay-3;
argL = kL(nlay)*R(nR);
argT = kT(nlay)*R(nR);
[JL,YL,JpL,YpL] = make_bessel(maxorder,argL,0);
[JT,YT,JpT,YpT] = make_bessel(maxorder,argT,0);
[F1L,F2L,G1L,G2L] = make_FG(maxorder,JL,YL,JpL,YpL,argL,0);
[F1T,F2T,G1T,G2T] = make_FG(maxorder,JT,YT,JpT,YpT,argT,0);
a1 = rho(nlay)*w^2;
a2 = 2*mu(nlay)*kL(nlay)/R(nR);
a3 = 2*muflag0(nlay)*rho(nlay)*w^2/(kT(nlay)*R(nR));
Z(irow,icol,:) = reshape(-a1*JL-a2*F1L,1,1,mmax);
Z(irow,icol+1,:) = reshape(-a1*YL-a2*F2L,1,1,mmax);
Z(irow,icol+2,:) = reshape(a3*m.*G1T,1,1,mmax);
Z(irow,icol+3,:) = reshape(a3*m.*G2T,1,1,mmax);
Z(irow+1,icol,:) = reshape(-a2*m.*G1L,1,1,mmax);
Z(irow+1,icol+1,:) = reshape(-a2*m.*G2L,1,1,mmax);
Z(irow+1,icol+2,:) = muflag0(nlay)*reshape(a1*JT+a3*F1T,1,1,mmax);
Z(irow+1,icol+3,:) = muflag0(nlay)*reshape(a1*YT+a3*F2T,1,1,mmax);
if R(1)==0
    Z(1:2,:,:)= 0;
    Z(:,[2 4],:) = 0;
end

% end function

```

```

function rout = cyl_outer_response(w,maxorder,layer_parms)
% function rout = cyl_outer_response(w,maxorder,layer_parms)
% This function calculates the frequency response for each mode up to
% maxorder for a layered cylinder with a unit force applied to the outside.
% Output is vector of modal amplitudes (displacement) on outer surface.
% Inputs are:
%   w:      angular frequency vector
%   maxorder: maximum order of angular dependence (0:maxorder)
%   layer_parms:    structure of layer parameters
%                 .R:      interface radii
%                 .mu:     shear modulus (Lame)
%                 .lambda: compression modulus (Lame)
%                 .rho:     density
% Created 2/23/07 from cyl_freq_response

nw = length(w);
mmax = maxorder+1;
rout = zeros(nw,mmax);
nlayers = length(layer_parms.mu);
src_vec = zeros(4*nlayers,1);
src_vec(4*nlayers-1,1) = 1;
for i=1:nw
    Z = make_Z(maxorder,layer_parms,w(i));
    [ampin,ampout] = measure_cyl(layer_parms,w(i),Z,maxorder,src_vec);
    rout(i,:) = ampout';
end

% end function

```

```

function [F1,F2,G1,G2] = make_FG(maxorder,J,Y,dJ,dY,x,zeroflag)
% function [F1,F2,G1,G2] = make_FG(maxorder,J,Y,dJ,dY,x,zeroflag)
% This function creates the F1, F2, G1, G2 functions from the
% Bessel functions and derivatives. If zeroflag is present set
% F2 and G2 to zero when x = 0.

m = (0:maxorder)';
msq = m.^2;
mlen = maxorder+1;
F1 = zeros(mlen,1);
F2 = zeros(mlen,1);
G1 = zeros(mlen,1);
G2 = zeros(mlen,1);
switch maxorder
    case 0
        if x==0
            if nargin==7
                Yovx = 0;
                Jovx = 0;
            else
                Yovx = Y/x;
                Jovx = J/x;
            end
        else
            mJovx = 0;
            mYovx = 0;
            Jovx = J/x;
            Yovx = Y/x;
        end
    case 1
        if x==0
            mJovx = [0 ; .5];
            if nargin==7
                mYovx = [0 ; 0];
                Jovx = [0 ; 0];
                Yovx = [0 ; 0];
            else
                mYovx = [0 ; Y/x];
                Jovx = J/x;
                Yovx = Y/x;
            end
        else
            mJovx = m.^2.*J/x;
            mYovx = m.^2.*Y/x;
            Jovx = J/x;
            Yovx = Y/x;
        end
    otherwise
        if x==0
            mJovx = [0 ; .5 ; zeros(mlen-2,1)];
            if nargin==7
                mYovx = zeros(mlen,1);
                Jovx = zeros(mlen,1);
                Yovx = zeros(mlen,1);
            else
                mYovx = [0 ; Y(2:mlen)/x];
                Jovx = J/x;
            end
        else
            mJovx = m.^2.*J/x;
            mYovx = m.^2.*Y/x;
            Jovx = J/x;
            Yovx = Y/x;
        end
    end
end

```

```

        Yovx = Y/x;
    end
else
    mJovx = m.^2.*J/x;
    mYovx = m.^2.*Y/x;
    Jovx = J/x;
    Yovx = Y/x;
end
end

F1 = dJ-mJovx;
F2 = dY-mYovx;
G1 = dJ-Jovx;
G2 = dY-Yovx;

% F1 = dJ-m.^2.*J/x;
% F2 = dY-m.^2.*Y/x;
% G1 = dJ-J/x;
% G2 = dY-Y/x;

% end function

```

```

function [Jn,Yn,dJn,dYn] = make_bessel(maxorder,arg,zeroflag)
% function [Jn,Yn,dJn,dYn] = make_bessel(maxorder,arg,zeroflag)
% This function calculates the Bessel function J, Y, and their
% derivatives (dJ, dY) for integer orders up to maxorder with a
% single argument. This is used to calculate acoustic fields for
% cylindrical objects. If zeroflag is present set Y and dY to
% zero when argument is zero.

nlen = maxorder+1;
Jn = besselj((0:(maxorder+1))',arg);
if nargin==3 & arg==0
    Yn = zeros(size(Jn));
else
    Yn = bessely((0:(maxorder+1))',arg);
end
dJn = zeros(nlen,1);
dYn = zeros(nlen,1);
dJn(1) = -Jn(2);
dYn(1) = -Yn(2);
dJn(2:nlen) = .5*(Jn(1:(nlen-1))-Jn(3:(nlen+1)));
dYn(2:nlen) = .5*(Yn(1:(nlen-1))-Yn(3:(nlen+1)));
Jn = Jn(1:nlen);
Yn = Yn(1:nlen);

% end function

```